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# An Electric-Circuit Model on the Inter-Tape Contact Resistance and Current Sharing for REBCO Cable and Magnet Applications

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**Abstract**—REBCO coated conductor has demonstrated high current capacity that can enable high-field magnets for high energy physics and fusion applications. However, quench protection is still one of the main challenges to be addressed for these applications. In addition,  $I_c$  and  $n$  value variations along the length of REBCO tapes exist in commercial production. The inter-tape contact resistance plays a key role to develop the self protection capability in cables and magnets by enabling current sharing and suppressing excessive eddy currents. Here we propose an electric-circuit model to describe the inter-tape contact resistance and its impact on the current sharing between REBCO tapes. We report the experiments on a 2-stacked tape REBCO cable with local  $I_c$  drop to validate the model. With the developed model, we study the upper limit of the contact resistance which allows current sharing between tapes. We also study the impact of variation in  $I_c$  and  $n$  values in tapes on the cable performance. Our model is expected to provide useful insight into the current sharing and target values for inter-tape contact resistance in REBCO cables and magnets for various applications.

**Index Terms**—Contact resistance, current sharing, REBCO, stacked tape cable.

## I. INTRODUCTION

REBCO ( $\text{REBa}_2\text{Cu}_3\text{O}_{7-\delta}$ , RE = rare earth) coated conductor is a promising superconductor with multiple potential applications given its superior performance in high-field and high current capacity. Different cable configurations of REBCO, such as twisted stacked tape cable (TSTC) [1], Conductor on Round Core (CORC®) [2] and Symmetric Tape Round REBCO (STAR) wires [3], have been proposed for next generation high-field magnets for accelerators and fusion applications. REBCO pancake coils and layer wound coils have also been applied in Nuclear Magnetic Resonance (NMR) and Magnetic Resonance Imaging (MRI) [4], [5].

Although REBCO conductors can enable diverse high-field magnet applications, the REBCO conductor and magnet technology need to address several significant challenges. The protection against quench under a high engineering current

density is an issue ( $100 - 1000 \text{ Amm}^{-2}$ ) [6]. Recent non-insulation (NI) REBCO magnet technology has demonstrated self-protection capability against quench and nominal operation regardless of defects by allowing current sharing between the conductor turns [7], [8], leading to a recent record DC field of 45.5 T [9]. Nevertheless, excessive eddy currents related to the low inter-tape contact resistance ( $R_c$ ) in NI coils cause magnetic field distortions and charging/discharging delays [10], [11].

Another challenge is that critical current ( $I_c$ ) and  $n$  value variations along the length of REBCO tapes exist in commercial tapes, which can cause non-uniform current distribution among the tapes in the cable, affecting the performance and reliable operation of cables and magnets [12]–[14].

The contact resistance between tapes plays a key role to determine the current distribution. Several models have been reported to understand the current distribution in REBCO cables. These models consider terminal resistances, tape inductances, and insulation between tapes [15]–[18]. It has been demonstrated that at steady-state conditions current sharing is determined by the terminal resistances for cables with insulation between tapes [19]–[22]. Different numerical methods and software have been implemented to analyze the electrical behavior of HTS cables [23]–[26].

In this paper, we report a simple electric-circuit model to analyze the impact of the contact resistance on current sharing between tapes with local  $I_c$  drops. We also study the impact of variation in  $I_c$  and  $n$  values on the transport performance of stacked tape cables. Our results can provide important insight into the role of inter-tape contact resistance on current sharing for REBCO cables and magnets.

## II. MODEL OF STACKED TAPE CABLE WITH CONTACT RESISTANCE

We consider a cable composed of  $N$  REBCO tapes. Each tape can be divided into  $M$  sections that are represented by a voltage source  $V_{i,j}$  (Fig. 1) following the power law,

$$V_{i,j} = V_c \left( \frac{I}{I_c} \right)^n, \quad (1)$$

where  $I_c$  is the critical current of the tape,  $n$  is the index value, and  $I$  the current flowing through the voltage source.  $V_c$  is the critical voltage  $LE_c$ , where  $L$  is the length of the tape and  $E_c$  is the electric field criterion of  $100 \mu\text{Vm}^{-1}$ . Both  $I_c$  and  $n$  value can vary along the tape and among the tapes.

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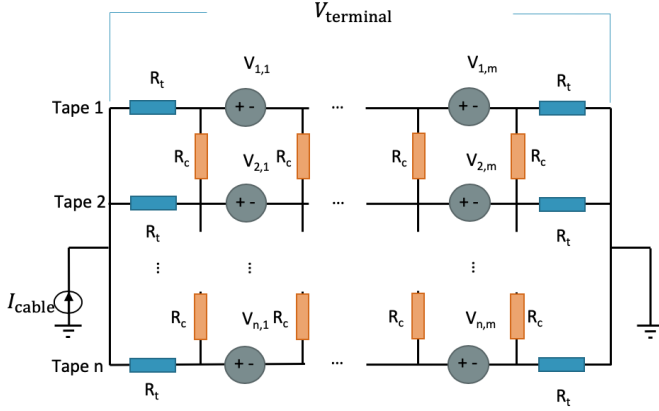


Fig. 1: A circuit model for stacked tape cable. REBCO tapes are divided into  $m$  sections represented by voltage sources  $V_{i,j}$ . Terminal resistances  $R_t$  and contact resistances  $R_c$  are included.

The model includes terminal resistances  $R_t$  at the ends of each tape, and contact resistances  $R_c$  between the tapes. All the tapes are connected to a current source  $I_{cable}$ . We use the open source NGSPICE [27], a circuit simulator based on the classic SPICE [28], to analyze the current and voltage distributions in the circuit model. The model can reproduce the results of measured and calculated current distributions on stacked tape cables [1], [15], [19].

### III. CURRENT SHARING IN A 2-STACKED TAPE CABLE WITH A LOCAL $I_c$ DROP

To study the impact of  $R_c$ , let us consider a simple 2-stacked tape cable with a local  $I_c$  drop in one of the tapes. We studied the voltage and current distributions with two cases: low and high  $R_c$ .

#### A. Experiments

We used two 4 mm wide SuperPower tapes (SCS4050) to make the cable. The tape has a 50  $\mu\text{m}$  thick Hastelloy® substrate and is surrounded by a 20  $\mu\text{m}$  thick Cu stabilizer. Tape 1 was 13 cm long and Tape 2 was 4 cm longer, such that we could solder both tapes directly to the current leads using Pb<sub>40</sub>Sn<sub>60</sub> solder. The substrate side of each tape was soldered directly to the terminals with 2 cm contact length at both ends. We soldered voltage taps separated by a distance of 2 cm in three sections on each REBCO tape. For Tape 1 these sections are labeled as V1, V2, and V3; for Tape 2 the corresponding sections are V4, V5, and V6 (Fig. 2). The voltage taps on Tape 1 were on the substrate side, whereas for Tape 2 they were on the REBCO side. The voltage taps were connected to nanovoltmeters to monitor the voltage. Another pair of voltage taps was soldered at the ends of Tape 2 to monitor the terminal voltage. We performed the following measurements:

- 1) Measure the  $I_c$  and  $n$  values for each tape separately at 77 K, self field.

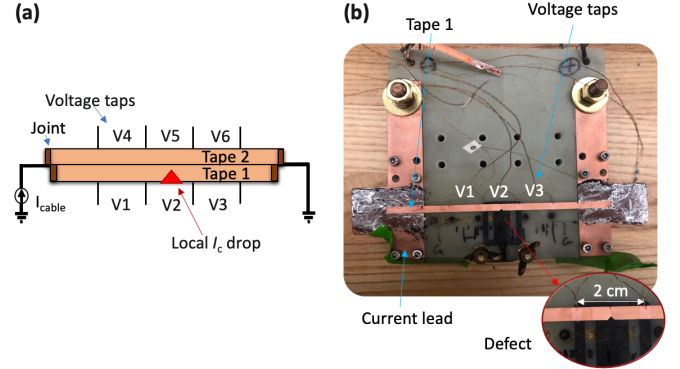


Fig. 2: 2-stacked tape cable with a local  $I_c$  drop. (a) Schematic of the cable (not to scale). (b) Experiment setup. We introduced the local  $I_c$  drop in the middle section of Tape 1 by making a small cut.

- 2) Create a defect in Tape 1 by cutting a small portion in section V2. Measure the reduced  $I_c$  and  $n$  value at 77 K, self field.
- 3) Stack the tapes with a Kapton tape in between for insulation (high  $R_c$ ). Solder the tapes to the current leads. Measure the voltage rise in all sections as a function of current.
- 4) Remove the Kapton tape and solder the two tapes together (low  $R_c$ ). Solder the tapes to the current leads while maintaining the same contact surface as in Step 3. Measure the voltage rise in all sections as a function of current.
- 5) Un-solder the tapes and repeat Step 1.

During the measurements, the current was ramped in a stair-step way and the voltage across each section was only measured when the current ramping stopped to minimize the inductive voltage. The measured  $I_c$  and  $n$  values are shown in Table I.  $I_c$  and  $n$  *before* are the values measured in Step 1 for all sections except V2, for which they were measured in Step 2.  $I_c$  and  $n$  *after* are the results from Step 5 after the tapes were un-soldered. Fig. 4(a) shows the measured  $V(I)$  data for the insulated case and Fig. 5(a) the  $V(I)$  data for the soldered case.

TABLE I: Measured  $I_c$  and  $n$  values for each tape section. Before being soldered into a stacked tape cable and after desoldering from the stacked tape cable. The values for V1 and V3 could not be measured after creating a defect in section V2.

Tape	Section	$I_c$ before (A)	$I_c$ after (A)	$n$ before	$n$ after
Tape 1	V1	132	—	30.5	—
Tape 1	V2	89	87	24.2	24.4
Tape 1	V3	133	—	31.0	—
Tape 2	V4	133	130	28.9	28.8
Tape 2	V5	130	130	30.7	29.9
Tape 2	V6	132	129	29.7	28.3

## B. Calculation

We applied the measured  $I_c$  and  $n$  values from Table I in the calculation. Fig. 3 shows the applied circuit model, which is a simplified case of the stacked tape cable model (Fig. 1) for two tapes with three sections each. The four contact resistances had the same constant value  $R_c$ . We did not consider the impact of the magnetic field on the  $I_c$  when both tapes are stacked. The terminal resistance  $R_t = 1.29 \text{ n}\Omega$  was determined from the measured terminal voltage before transition. We set  $R_c = 1000 \Omega$  for the insulated case. For the soldered case,  $R_c$  was  $46 \text{ n}\Omega$  based on the conductivity of the solder layer of  $0.1 \text{ mm}$  thickness ( $27 \text{ MSm}^{-1}$  at  $77 \text{ K}$  [29]). The calculated voltages at the tape sections for high and low  $R_c$  are shown in Fig. 4(b) and Fig. 5(b). The current distribution through the tape sections and contact resistances are shown in Figs. 6 and 7. For the case of low  $R_c$ , the path of these currents is shown in Fig. 3. The power generated at the cable terminals and the section V2 for high and low  $R_c$  are shown in Fig. 8.

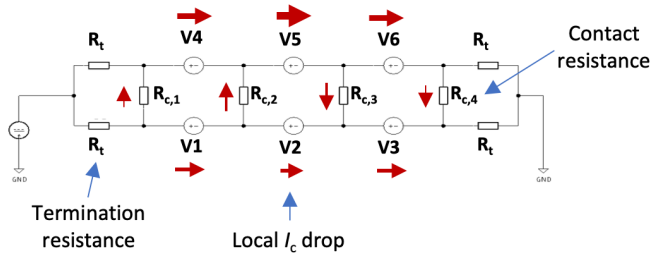


Fig. 3: Circuit model of a 2-stacked tape cable. The arrows show the current path for the case of low  $R_c$ . The size of the arrows indicates the amount of current (not to scale).

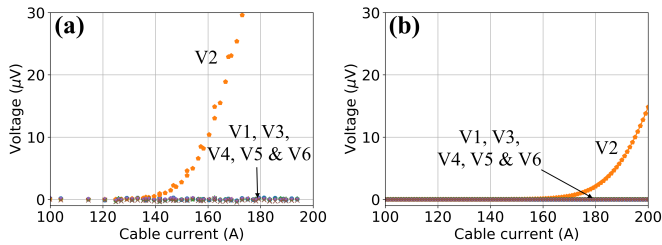


Fig. 4: Voltage at the tape sections with insulation (high  $R_c$ ). (a) Measurement. (b) Calculation.

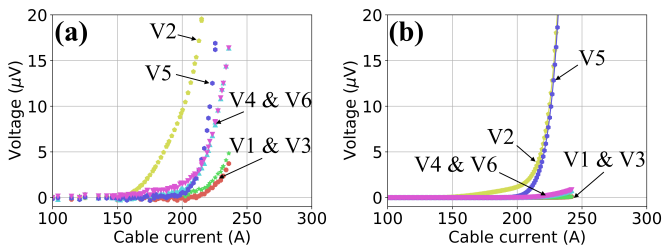


Fig. 5: Voltage at the tape sections with solder (low  $R_c$ ). (a) Measurement. (b) Calculation.

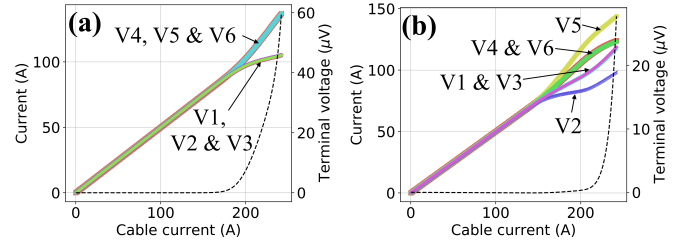


Fig. 6: Current through the tape sections. The dashed line is the terminal voltage with the resistive foot removed. (a) High  $R_c$ . (b) Low  $R_c$ .

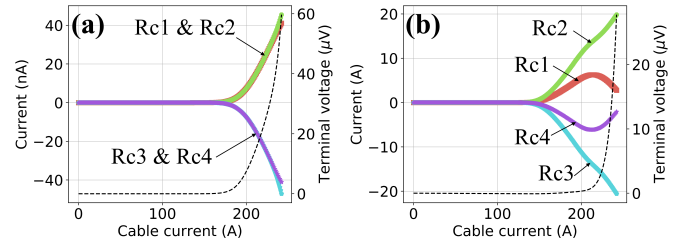


Fig. 7: Current distribution through the contact resistances. The dashed line is the terminal voltage with the resistive foot removed. (a) High  $R_c$ . (b) Low  $R_c$ .

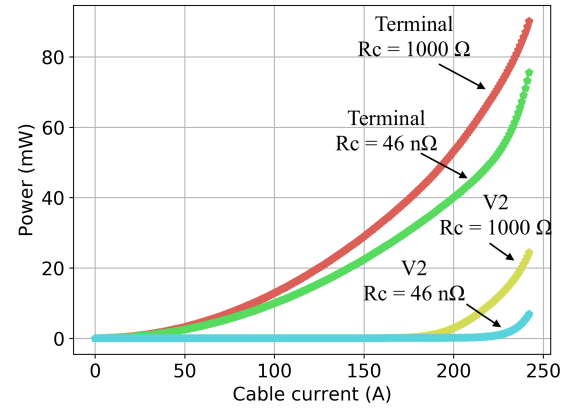


Fig. 8: Power generated at the cable terminals and the section V2 for high  $R_c = 1000 \Omega$  and low  $R_c = 46 \text{ n}\Omega$ .

## IV. UPPER LIMIT OF $R_c$ TO ALLOW CURRENT SHARING

The model presented here can also be used to determine the upper limit of  $R_c$  that allows current sharing. As an example, let us consider the 2-stacked tape cable model (Fig. 3). We assumed uniform  $I_c$  in each tape, having a value of  $50 \text{ A}$  in Tape 1 and  $100 \text{ A}$  in Tape 2. The  $n$  value was constant in all sections and equal to 30. The contact resistances  $R_{c,i}$  were also uniform. We swept the  $R_c$  from  $1 \text{ p}\Omega$  to  $1000 \Omega$  for three  $R_t$  values:  $2.5 \text{ n}\Omega$ ,  $250 \text{ n}\Omega$ , and  $25 \mu\Omega$ . We calculated the terminal voltage and the current distribution through the contact resistances as a function of current. Fig. 9 shows the results for  $R_t = 250 \text{ n}\Omega$ . The upper limit of  $R_c$  was selected in such a way that the terminal voltages and current distributions

for  $R_c$  values below this limit showed similar results. The upper limits of  $R_c$  to allow current sharing depending on  $R_t$  are shown in Table II.

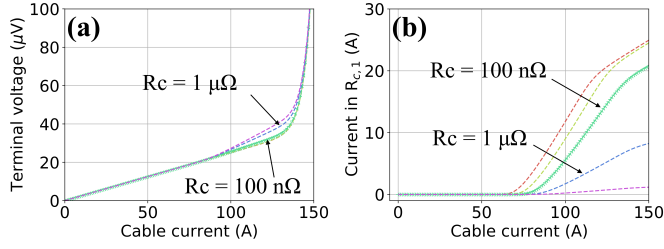


Fig. 9: Calculation results for various  $R_c$  in a 2-stacked tape cable with  $R_t = 250 \text{ n}\Omega$ . (a) Terminal voltage for different  $R_c$ . Voltages for  $R_c < 100 \text{ n}\Omega$  are overlapped. (b) Current through the contact resistance  $R_{c,1}$ .

TABLE II: Calculated upper limit of  $R_c$  to allow current sharing depending on  $R_t$  in a 2-stacked tape cable.

$R_t$	$R_c$ upper limit
2.5 nΩ	1 nΩ
250 nΩ	100 nΩ
25 μΩ	10 μΩ

## V. IMPACT OF $I_c$ AND $n$ VALUE VARIATION AMONG THE TAPES OF THE STACKED TAPE CABLE

We performed Monte Carlo simulations with sample size of 500 to study the impact of  $I_c$  and  $n$  value variation in the stacked tape cable model. The number of tapes ranges from 10 to 40. We analyzed two cases: high  $R_c = \infty$  (open circuit) and low  $R_c = 10 \text{ n}\Omega$ . Tape  $I_c$  and  $n$  values were generated by normal distributions with a mean  $I_c$  of 100 A and a mean  $n$  value of 30. The standard deviations of  $I_c$  and  $n$  values were varied up to 50% of their mean values.

The terminal voltage  $V_{\text{terminal}}$  was calculated by sweeping the current up to 150% the cable  $I_c$  and time dependence was not considered.  $V_{\text{terminal}}$  was limited to 100 μV and fitted using (1) with  $E_c = 100 \mu\text{Vm}^{-1}$  and  $L = 1 \text{ m}$  to calculate the cable  $I_c$  and  $n$  value.

The cable  $I_c$  and  $n$  values for the 500 samples per each standard deviation were averaged and normalized to the results of 1% standard deviation. Figs. 10 and 11 show the results from the variations of  $I_c$  and  $n$  value.

## VI. DISCUSSION

Calculated voltage across each tape section qualitatively agreed with measurements in both insulation and solder cases. For high  $R_c$  only voltage in section V2 rose (Fig. 4). In contrast, for low  $R_c$  all sections showed voltage rise (Fig. 5), which is an indication of current sharing. There are several factors that can explain the discrepancy between the measurements and calculations: the actual values of  $R_c$  and  $R_t$  may be different from the values used in the calculation; the  $I_c$  reduction due to self-field effects in the stacked tape cable

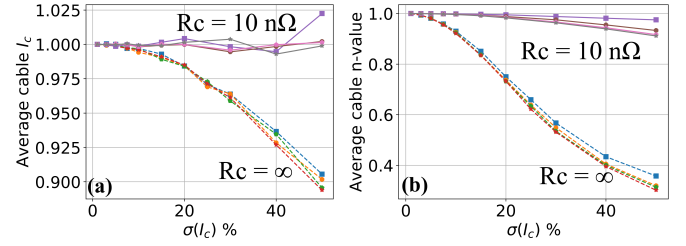


Fig. 10: Average cable parameters as a function of  $I_c$  variation for different number of tapes.  $n = 30$  for all tapes. (a) Cable  $I_c$ . (b) Cable  $n$  value.

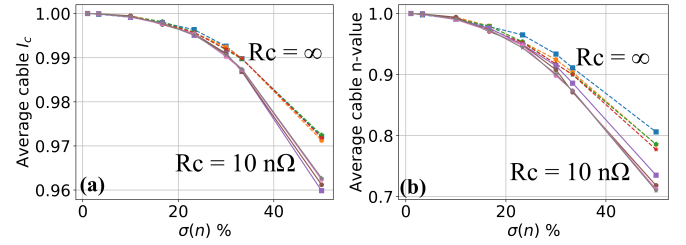


Fig. 11: Average cable parameters as a function of  $n$  value variation for different number of tapes.  $I_c = 100 \text{ A}$  for all tapes. (a) Cable  $I_c$ . (b) Cable  $n$  value.

[30] was not considered; potential non-uniform  $I_c$  within each section can also contribute to the discrepancy. Although the  $I_c$  from each tape decreased by 3% after being soldered, the voltage rise in the soldered case (Fig. 5 (a)) was not caused by degradation of the tapes.

Low contact resistance allowed current to bypass the section with local  $I_c$  drop when approaching to its critical current (Figs. 3, 6(b), and 7(b)). Low  $R_c$  is necessary to decrease the power generation and the resulting temperature rise during the cable transition to normal state. For instance, with the low  $R_c$  the power was reduced by 26% at the cable terminals and by 90% at the section V2 when the cable current was 225 A (Fig. 8).

One important question is how low the contact resistance should be to allow current sharing. Our model can address this question by providing insight on the current distribution with different values of  $R_c$  (Fig. 9). For a simple 2-stacked tape cable, our model suggests that the  $R_c$  should be lower than 40% of the terminal resistance.

The simple circuit model presented here also allowed us to assess how the variation of  $I_c$  and  $n$  values in individual tapes affects the  $I_c$  and  $n$  value of the cable. With high  $R_c$  the cable  $I_c$  is less affected than the cable  $n$  value when the tape  $I_c$  or  $n$  value varies (Figs. 10 and 11). For example, a standard deviation of  $I_c$  of 30% only leads to 4% of cable  $I_c$  reduction, but it causes 50% drop in  $n$  value (Fig. 10). This result may also be interpreted as the following: if a cable with a high  $R_c$  between tapes shows a low  $n$  value (e.g.  $< 15$ , 50% of the mean value used in the model) this may indicate a strong  $I_c$  variation among the tapes.



We found that the cable  $I_c$  and  $n$  value were not affected by  $I_c$  variation when current can share between tapes through low  $R_c$  (Fig. 10). Hence, a low  $R_c$  between tapes benefits the cable performance by allowing current sharing, which is consistent with the observation in [8].

## VII. CONCLUSION

We developed a simple circuit model to help understand the impact of contact resistance and current sharing for REBCO cables. The model can provide important insight of the relevance of the inter-tape contact resistance in REBCO cables. In particular, the benefits that low  $R_c$  brings to cable performance in the presence of defects.

The calculation qualitatively agreed with the measurements in a 2-stacked tape cable with solder or insulation between tapes. The current bypassed the section with lower  $I_c$  when the contact resistance was low. As a result, the power generation was lower compared to the case of high contact resistance. The model also gave insight on the current distribution for different  $R_c$  values. Reducing  $R_c$  will help the cable tolerate the variation of  $I_c$  in the tapes. We intend to include the self and mutual inductances between tapes to study the current sharing during non-steady state conditions.

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